

Phys 5701 13 Oct 2020

Today: "Choice of Gauge" in Quantum Mechanics

Final class on "Quantum Dynamics"

Thursday will start "Angular Momentum"

Reminder from Classical Physics

$$V(\vec{x}) \rightarrow \tilde{V}(\vec{x}) \equiv V(\vec{x}) + V_0$$

$$m\vec{a} = \vec{F} = -\vec{\nabla}V = -\vec{\nabla}\tilde{V}$$

No change in physics!

Big Effect in Quantum Mechanics!

$$|\alpha, t_0; t\rangle = e^{-iH(t-t_0)/\hbar} |\alpha, t_0\rangle$$

$$\rightarrow \widetilde{|\alpha, t_0; t\rangle} = e^{-i\tilde{H}(t-t_0)/\hbar} |\alpha, t_0\rangle$$

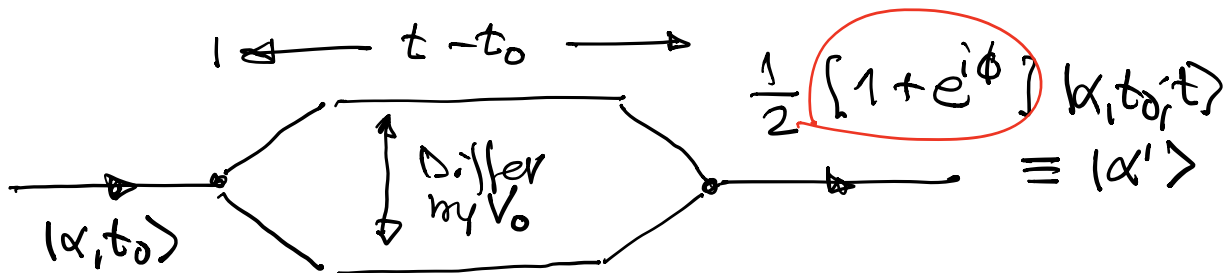
$$= e^{-iV_0(t-t_0)/\hbar} |\alpha, t_0; t\rangle$$

$$= e^{-i\phi(t)/\hbar} \widetilde{|\alpha, t_0; t\rangle}$$

But $\langle \beta | \tilde{\alpha}, t_0; t \rangle = e^{-i\phi/\hbar} \langle \beta | \alpha, t_0; t \rangle$

$\Leftrightarrow |\langle \beta | \tilde{\alpha}, t_0; t \rangle|^2 = |\langle \beta | \alpha, t_0; t \rangle|^2$

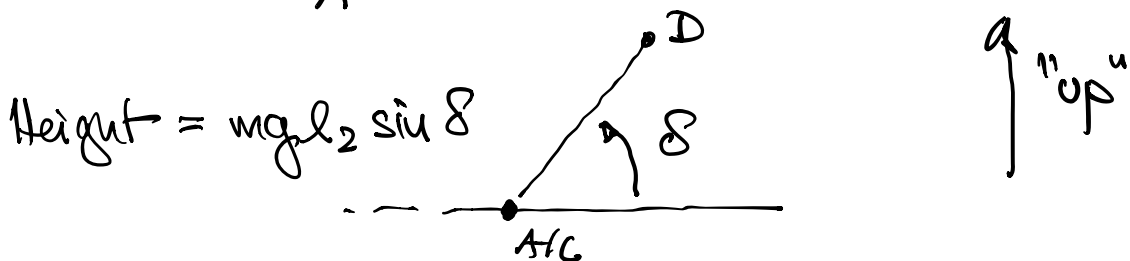
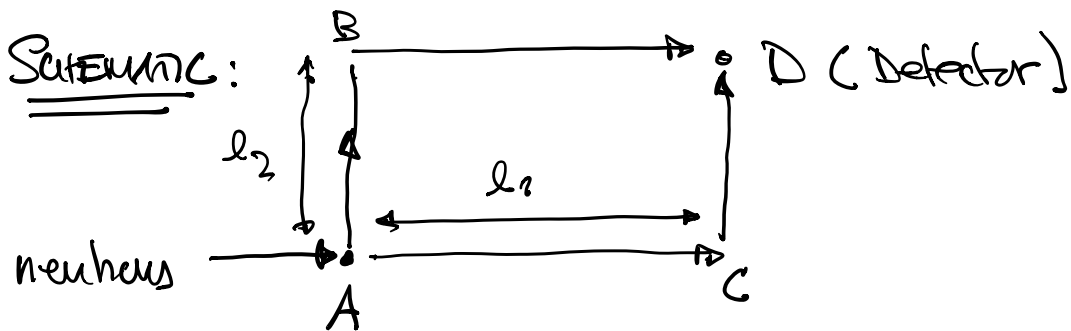
However consider...



$\Leftrightarrow |\langle \beta | \alpha' \rangle|^2 = \frac{1}{4} |1 + e^{i\phi}|^2 |\langle \beta | \alpha, t_0; t \rangle|^2$

i.e. Probability oscillates as $\frac{1}{2} (1 + \cos \phi(t))$
 with $\phi(t) = V_0(t-t_0)/\hbar$

SEE PRL 34 (1975) 1472 **POSTED!**



Review of Electromagnetism App. A, C

Maxwell's Equations in CGS UNITS!

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho \quad \vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \cdot \vec{B} = 0 \quad \vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

Lorentz Force Law

$$\vec{F} = q\vec{E} + \frac{1}{c} q\vec{v} \times \vec{B}$$

$\vec{v} = \dot{\vec{x}}$

Potentials (i.e. "electromagnetic potentials")

$$\vec{B}(\vec{x}, t) = \vec{\nabla} \times \vec{A}(\vec{x}, t)$$

$$\vec{E}(\vec{x}, t) = -\vec{\nabla}\phi(\vec{x}, t) - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

1e. Discuss \vec{A} and ϕ not \vec{E} and \vec{B}

STATIC

Lagrangian

$$L(\vec{x}, \dot{\vec{x}}) = \frac{1}{2} m \dot{\vec{x}}^2 - q\phi(\vec{x}) + \frac{q}{c} \dot{\vec{x}} \cdot \vec{A}(\vec{x})$$

Canonical Momentum

$$\vec{p} = \partial L / \partial \dot{\vec{x}} = m \dot{\vec{x}} + \frac{q}{c} \vec{A}(\vec{x})$$

Hamiltonian

$$\begin{aligned} H(\vec{x}, \vec{p}) &\equiv \vec{x} \cdot \vec{p} - L(\vec{x}, \dot{\vec{x}}(\vec{p})) \\ &= \frac{1}{2m} \left[\vec{p} - \frac{q}{c} \vec{A}(\vec{x}) \right]^2 + q\phi(\vec{x}) \end{aligned}$$

Electric Force: $q\phi(\vec{x}) = U(\vec{x})$

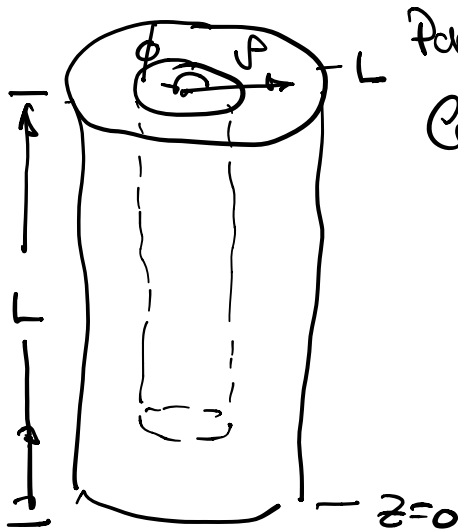
↳ Same as for "gravity"

Magnetic Field: More interesting!

$$\vec{A}(\vec{x}) \rightarrow \vec{\tilde{A}}(\vec{x}) = \vec{A}(\vec{x}) + \vec{\nabla} \Lambda(\vec{x})$$

$$\begin{aligned} \text{↳ } \vec{B} &= \vec{\nabla} \times \vec{\tilde{A}} = \underbrace{\vec{\nabla} \times \vec{A}}_{\vec{B}} + \underbrace{\vec{\nabla} \times \vec{\nabla} \Lambda}_{\equiv 0} \end{aligned}$$

Flux Quantization Prob. 2.37



Particle mass μ

Cylindrical coordinates
 (ρ, ϕ, z)

Particle confined to

$$\underline{\underline{\rho_a \leq \rho \leq \rho_b}}$$

"Free Particle"

$$-\frac{\hbar^2}{2\mu} \nabla^2 \psi(\rho, \phi, z) = E \psi(\rho, \phi, z)$$

$$\nabla = \hat{\rho} \frac{\partial}{\partial \rho} + \hat{z} \frac{\partial}{\partial z} + \hat{\phi} \frac{1}{\rho} \frac{\partial}{\partial \phi}$$

$$\nabla^2 = \nabla \cdot \nabla$$

$$\text{L.H.S. } \nabla^2 \psi = \nabla \cdot \nabla \psi$$

$$= \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \psi}{\partial \rho} \right) + \frac{\partial^2 \psi}{\partial z^2} + \frac{1}{\rho^2} \frac{\partial^2 \psi}{\partial \phi^2}$$

• Separate variables: $\psi = R(\rho) \Phi(\phi) Z(z)$

• Boundary conditions

$$R(\rho_a) = R(\rho_b) = 0$$

$$Z(0) = 0 = Z(L)$$

$$\Phi(\phi + 2\pi) = \Phi(\phi)$$

• Write $E = \frac{\hbar^2 k^2}{2\mu}$

For homework, show that

$$\Phi(\phi) = e^{\pm im\phi} \quad m = 0, 1, 2, 3, \dots$$

$$Z(z) = e^{\pm i\alpha_l z} \quad \alpha_l \equiv \frac{\pi}{L} l \quad l = 1, 2, 3, \dots$$

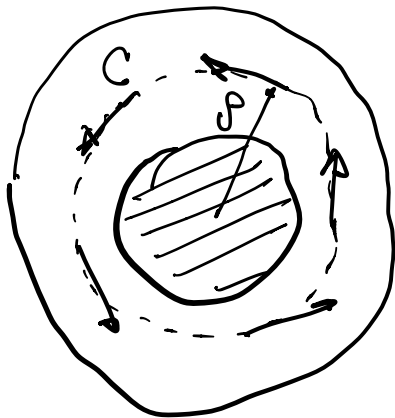
$$\frac{1}{\rho} \frac{d}{d\rho} \left(\rho \frac{dR}{d\rho} \right) + \left(k^2 - \alpha_l^2 - \frac{m^2}{\rho^2} \right) R(\rho) = 0$$

\Leftrightarrow Bessel's Equation for $x = k\rho$

$$k^2 \equiv k^2 - \alpha_l^2$$

then "Eigenvalues" for energy.

Now add \vec{B} field: $\rho \leq \rho_a$



$$\oint_C \vec{A} \cdot d\vec{\ell} = A (2\pi\rho)$$

$$= \int_S (\vec{\nabla} \times \vec{A}) \cdot d\vec{S}$$

$$= \int_S \vec{B} \cdot d\vec{S} = B \pi \rho_a^2$$

$$\Rightarrow \vec{A} = \vec{A}(\rho) = \frac{B \rho_a^2}{2\rho} \hat{\phi}$$

"this affects things"

$$\vec{p} \rightarrow \vec{p} - \frac{e}{c} \vec{A} \quad q = e < 0$$

$$\text{i.e. } \frac{\partial}{\partial t} \rightarrow \frac{\partial}{\partial t} - \frac{ie}{\hbar c} \frac{B \rho_a^2}{2}$$

Using Gauge Transformations

$$\text{i.e. } \vec{A} \rightarrow \vec{\tilde{A}} = \vec{A} + \vec{\nabla}\Lambda = 0$$

$$\text{where } \Lambda(\phi) = - \frac{B \rho_a^2 \phi}{2}$$

$$\underline{\text{Book}}: \vec{A} \rightarrow \vec{\tilde{A}} \Rightarrow |\alpha\rangle \rightarrow |\tilde{\alpha}\rangle = G |\alpha\rangle$$

$$G^\dagger G = 1$$

$$\text{want } \langle \tilde{\alpha} | \left[\vec{p} - \frac{e \vec{A}^*}{c} \right] | \alpha \rangle = \langle \tilde{\alpha} | \left[\vec{p} - \frac{e \vec{A}}{c} \right] | \tilde{\alpha} \rangle$$

$$\Leftrightarrow G = G(\vec{x}) = \exp \left[\frac{ie}{\hbar c} \Lambda(\vec{x}) \right]$$

For our problem

$$G = G(\phi) = e^{-ig\phi}$$

$$g = eB\rho_a^2 / 2\hbar c$$

$$\Leftrightarrow G(\phi) = e^{ig\phi} \quad \gamma = m \pm g$$

$$\Rightarrow g = \pm m \Rightarrow \boxed{B \pi \rho_a^2 = \pm \frac{\hbar c}{e} m} \quad 4 \times 10^{-7} \text{ gauss} \cdot \text{cm}^2$$