

Phys 4702 Tuho QM II Fall 2024 17 Oct 2024

NOTE: Re $\langle l^2 | l^2 | l^2 | l^2 \rangle$

* HW #8 status

$$[A, B] = 0$$

$\Rightarrow \langle a | B | a \rangle = \langle a |$ + office hours tomorrow (I have to leave after!)

Today: Discrete Translation Symmetry

Where we are going: Bloch's theorem

\hookrightarrow Crystal Energy Bands / Gaps (HW!)

Recall: $T(\Delta x') |x'\rangle = |x' + \Delta x'\rangle$ [Easy to generalize

to 2D or 3D!]

$$\hookrightarrow T(\Delta x') = \exp\left[-\frac{i p \Delta x'}{\hbar}\right]$$

p = momentum (in x -direction); unitary operator

Detail: Eigenvalues of Unitary Operators

[Not Hermitian \Rightarrow Not observable, OK, just go!]

$$U|\lambda\rangle = \lambda|\lambda\rangle \quad \text{w/} \quad U^\dagger = U^{-1} \quad \text{i.e.} \quad U^\dagger U = 1$$

$$\hookrightarrow \langle \lambda | U^\dagger = \lambda^* \langle \lambda | \Rightarrow (\langle \lambda | U^\dagger)(U | \lambda \rangle) = \lambda^* \lambda \langle \lambda | \lambda \rangle$$

$$\text{But } U^\dagger U = 1 \Rightarrow \langle \lambda | (U^\dagger U) | \lambda \rangle = \langle \lambda | \lambda \rangle = |\lambda|^2 \langle \lambda | \lambda \rangle$$

i.e. $|\lambda|^2 = 1$ [Recall "positivity postulate" i.e. $\langle \lambda | \lambda \rangle > 0$]

\hookrightarrow Eigenvalue of unitary operator $\lambda = e^{i\alpha}$, $\alpha \in \mathbb{R}$

Note: Parity is a special case (Hermitian!)

↳ Eigenvalues must be real \Rightarrow only ± 1

Discrete Translation...

$$T(a)|x'\rangle = |x'+a\rangle \quad \text{w/ } T(a) = \exp\left(-\frac{i}{\hbar} p a\right)$$

$$x T(a)|x'\rangle = x|x'+a\rangle = (x'+a)|x'+a\rangle$$

$$T^\dagger(a) x T(a)|x'\rangle = (x'+a) T^\dagger|x'+a\rangle$$

$$\text{But } T^\dagger(a) = \exp\left(+\frac{i}{\hbar} p a\right) = \exp\left(-\frac{i}{\hbar} p(-a)\right)$$

$$= T(-a) \quad [\text{of course! Time is reverse!}]$$

$$\text{↳ } T^\dagger(a) x T(a)|x'\rangle = (x'+a)|x'\rangle = (x+a)|x'\rangle$$

i.e. $T^\dagger(a) x T(a) = x + a$ Makes sense! Good.

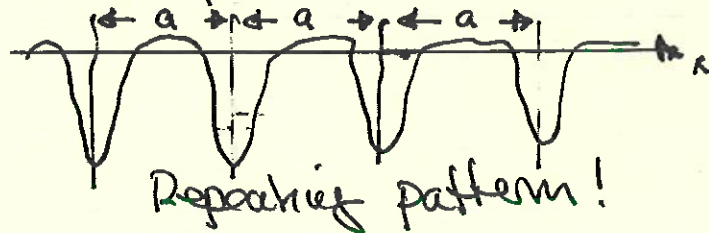
Note: $T^\dagger(a) F(p) T(a) = F(p)$ (i.e. $[T(a), p] = 0$)

... Symmetry: Require $T^\dagger(a) H T(a) = H$
 $\text{↳ } [H, T(a)] = 0$

$$H = \frac{1}{2m} p^2 + V(x)$$

Symmetric!!

Need "periodicity" i.e.



$$\text{i.e. } T^\dagger(a) V(x) T(a) = V(x+a) = V(x)$$

$$\Leftrightarrow \underline{T^\dagger(a) H T(a)} = H \quad \text{for } V(x) \text{ w/ period } a$$

"Discrete Translation Symmetry"

Eigenstates

$$\text{Translation symmetry} \Rightarrow [H, T(a)] = 0$$

"Simultaneous Eigenvalues for H and $T(a)$ "

We know eigenvalues of $T(a)$ are $e^{i\alpha}$

\Leftrightarrow How are these related to eigenvalues E of H ?

Work in position space:

$$T(a) |\alpha\rangle = e^{i\alpha} |\alpha\rangle$$

$$\langle x | T(a) |\alpha\rangle = e^{i\alpha} \langle x | \alpha\rangle$$

$$\langle x-a | \alpha\rangle = e^{i\alpha} \langle x | \alpha\rangle$$

Define $\psi(x) \equiv \langle x | \alpha\rangle$

$$\Leftrightarrow \psi(x-a) = e^{i\alpha} \psi(x)$$

Convention: $\alpha = -ka$ NOTE: $\hbar k$ has units of momentum.

$$\Leftrightarrow \psi(x-a) = e^{-ika} \psi(x)$$

Now write $\psi(x) = e^{ikx} u(x)$ for some $u(x)$

$$\Leftrightarrow e^{ik(x-a)} u(x-a) = e^{-ika} e^{ikx} u(x)$$

$$u(x-a) = u(x)$$

i.e. $\psi(x) = e^{ikx} u(x)$ w/ $u(x+ta) = u(x)$
"Periodic"

"Bloch's Theorem"

... Now explore the energy eigenvalues!

NOTE: e^{ikx} looks like a free particle wave of $p = \hbar k$

NOTE: $\psi(x+ta) = e^{ik(x+ta)} u(x+ta)$
 $= e^{ikta} e^{ikx} u(x) = e^{ikta} \psi(x)$ ← Another way to state Bloch's Theorem

$\implies |\psi(x+ta)|^2 = |\psi(x)|^2$ Makes sense!

Energy Eigenvalues

Substitute $\psi(x) = e^{ikx} u(x)$ into Schrödinger Equ.

\implies Expect to find $E = E(k)$.

What can we say about the values of k ?

First consider "infinite" periodicity as $N \sim 10^{23}$ sites that wrap around start to end.

i.e. $\psi(x+Na) = \psi(x) \implies \exp(Nika) = 1$

$\implies k = \frac{2\pi}{Na} n \quad n = 0, \pm 1, \pm 2, \dots$

But N is very large!

\implies Take k to be a continuous variable.

For specific problem need to know $V(x)$!

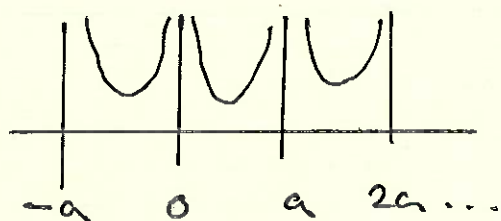
Homework: Periodic δ -functions

$$V(x) = \frac{\hbar^2}{2m} \lambda \sum_{n=-\infty}^{\infty} \delta(x-na)$$

wave function is just "free particle" for $0 < x < a$
 [the problem guides you \Rightarrow "Energy Bands"]

General Considerations

1) Infinite Barrier: completely localized state $|n\rangle$



$$H|n\rangle = \underline{\underline{E_0(n)}} \text{ for all } n \quad (\text{Ground states})$$

"Infinite Degeneracy"

But $T(a)|n\rangle = |n+1\rangle \Rightarrow |n\rangle$ not eigenstates of $T(a)$!

Try the linear combination

$$|\theta\rangle = \sum_{n=-\infty}^{\infty} e^{in\theta} |n\rangle \quad \text{SHOULD } H|\theta\rangle = E_0|\theta\rangle$$

$$\begin{aligned} \text{But } T(a)|\theta\rangle &= \sum_{n=-\infty}^{\infty} e^{in\theta} |n+1\rangle = \sum_{n=-\infty}^{\infty} e^{i(n-1)\theta} |n\rangle \\ &= e^{-i\theta} |\theta\rangle \quad \text{Simultaneous Eigenstates!} \end{aligned}$$

NOTE:

$$(2) \langle n-1 | H | n \rangle = \int dx \langle n-1 | x \rangle \langle x | H | n \rangle$$

$$(1) \langle x-a | \theta \rangle = e^{-i\theta} \langle x | \theta \rangle$$

$$\text{i.e. } \underline{\underline{\theta = \alpha = ka}}$$

$$\frac{\hbar^2}{2m} \psi_{n-1}(x) \left[-\frac{\hbar^2}{2m} \frac{d^2 \psi_n}{dx^2} + V(x) \psi_n(x) \right]$$

= 0 (NO OVERLAP DUE TO INFINITE BARRIER)

2) Finite (But "high") Barrier $\Rightarrow \langle n-1 | H | n \rangle \neq 0$
 "off diagonal elements" so must diagonalize
 for energy eigenvalues!

write diagonal elements as $\langle n | H | n \rangle = E_0$

assume $\langle n' | H | n \rangle = 0$ unless $n' = n \pm 1$

"Tight Binding Approximation"

Let $\langle n \pm 1 | H | n \rangle = -\Delta$

Also $\langle n \pm 1 | n \rangle = \int dx \langle n \pm 1 | x \rangle \langle x | n \rangle \approx 0$

$$\Rightarrow H | n \rangle = E_0 | n \rangle - \Delta | n-1 \rangle - \Delta | n+1 \rangle \quad \left[= \sum_{m=-\infty}^{\infty} \langle m | H | n \rangle | m \rangle \right]$$

But is $|\theta\rangle$ still an energy eigenstate?

$$\begin{aligned} H |\theta\rangle &= H \sum_{n=-\infty}^{\infty} e^{in\theta} | n \rangle = \sum_{n=-\infty}^{\infty} e^{in\theta} [E_0 | n \rangle - \Delta | n-1 \rangle - \Delta | n+1 \rangle] \\ &= \sum_{n=-\infty}^{\infty} e^{in\theta} E_0 | n \rangle - \sum_{n=-\infty}^{\infty} \Delta [e^{i(n+1)\theta} + e^{i(n-1)\theta}] | n \rangle \\ &= [E_0 - \Delta (e^{i\theta} + e^{-i\theta})] \sum_{n=-\infty}^{\infty} e^{in\theta} | n \rangle \end{aligned}$$

$\Rightarrow H |\theta\rangle = (E_0 - 2\Delta \cos \theta) |\theta\rangle$ yes! Eigenstate!

$$E = E_0 - 2\Delta \cos \theta = E_0 - 2\Delta \underline{\underline{\cos(ka)}}$$

Between -1 and 1

i.e. "Energy Band" w $E_0 - 2\Delta \leq E \leq E_0 + 2\Delta$

[Remember: $\hbar\omega$ is a specific example]