

[This class over Zoom. Fingers crossed!]

Review so far

Hermitian Operators and Basis States

$$A|a'\rangle = a'|a'\rangle \Rightarrow 1 = \sum_{a'} |a'\rangle \langle a'|$$

$$\Leftrightarrow |a\rangle = 1|a\rangle = \sum_{a'} |a'\rangle \langle a'|a\rangle \text{ coefficients}$$

Spin-1/2 Example:

$$S_z | \pm \hat{z} \rangle = \pm \frac{\hbar}{2} | \pm \hat{z} \rangle$$

$$| \pm \hat{x} \rangle = \frac{1}{\sqrt{2}} | + \hat{z} \rangle \pm \frac{1}{\sqrt{2}} | - \hat{z} \rangle$$

$$| \pm \hat{y} \rangle = \frac{1}{\sqrt{2}} | + \hat{z} \rangle \pm \frac{i}{\sqrt{2}} | - \hat{z} \rangle$$

Building Operators* From Outer Products

$$A = A1 = A \sum_{a'} |a'\rangle \langle a'| = \sum_{a'} a' |a'\rangle \langle a'|$$

* For now, Hermitian Operators

Spin-1/2 Example:

$$S_z = \frac{\hbar}{2} [| + \hat{z} \rangle \langle + \hat{z} | - | - \hat{z} \rangle \langle - \hat{z} |]$$

$$S_x = \frac{\hbar}{2} [| + \hat{z} \rangle \langle - \hat{z} | + | - \hat{z} \rangle \langle + \hat{z} |]$$

$$S_y = \frac{\hbar}{2} [-i | + \hat{z} \rangle \langle - \hat{z} | + i | - \hat{z} \rangle \langle + \hat{z} |]$$

Today: Matrix Representations

It's how we actually solve many QM problems.

Key point: Pick a basis

↳ kets, bras, operators are represented in that basis

Different basis, Different values, but same operators/states!!

Simple Example: Spin-1/2 in the $| \pm \hat{z} \rangle$ Basis

$$\begin{aligned} | \alpha \rangle &= \frac{1}{\sqrt{2}} | \alpha \rangle = (| + \hat{z} \rangle \langle + \hat{z} | + | - \hat{z} \rangle \langle - \hat{z} |) | \alpha \rangle \\ &= | + \hat{z} \rangle \underbrace{\langle + \hat{z} | \alpha \rangle} + | - \hat{z} \rangle \underbrace{\langle - \hat{z} | \alpha \rangle} \end{aligned}$$

↳ the two complex numbers $\langle + \hat{z} | \alpha \rangle$ and $\langle - \hat{z} | \alpha \rangle$ completely specify $| \alpha \rangle$ if we agree to use $| \pm \hat{z} \rangle$ basis!

It is handy to write this representation as

$$| \alpha \rangle \doteq \underline{\alpha} = \begin{bmatrix} \langle + \hat{z} | \alpha \rangle \\ \langle - \hat{z} | \alpha \rangle \end{bmatrix} \quad \text{"column vector"}$$

↑ Peculiar symbol but this is NOT "="!

$$\text{e.g. } | + \hat{z} \rangle \doteq \begin{bmatrix} 1 \\ 0 \end{bmatrix} \equiv \chi_+ \quad (\text{in some books})$$

$$| - \hat{z} \rangle \doteq \begin{bmatrix} 0 \\ 1 \end{bmatrix} \equiv \chi_-$$

$$| + \hat{y} \rangle = \frac{1}{\sqrt{2}} | + \hat{z} \rangle + \frac{i}{\sqrt{2}} | - \hat{z} \rangle \doteq \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix} \quad \text{etc...}$$

Similar for bras \Rightarrow Row (Dual) vectors

$$\begin{aligned} \langle \alpha | &= \langle \alpha | 1 = \langle \alpha | +\hat{z} \rangle \langle +\hat{z} | + \langle \alpha | -\hat{z} \rangle \langle -\hat{z} | \\ &= \langle +\hat{z} | \alpha \rangle^* \langle +\hat{z} | + \langle -\hat{z} | \alpha \rangle^* \langle -\hat{z} | \\ &\doteq \underline{\tilde{\alpha}} = [\langle +\hat{z} | \alpha \rangle^* \quad \langle -\hat{z} | \alpha \rangle^*] \quad \text{etc...} \end{aligned}$$

In General

$$|\alpha\rangle = 1|\alpha\rangle = \sum_{a'} |a'\rangle \langle a' | \alpha \rangle \doteq \begin{bmatrix} \langle a^{(1)} | \alpha \rangle \\ \langle a^{(2)} | \alpha \rangle \\ \vdots \\ \langle a^{(N)} | \alpha \rangle \end{bmatrix}$$

$$\langle \alpha | = [\langle a^{(1)} | \alpha \rangle^* \quad \langle a^{(2)} | \alpha \rangle^* \quad \dots \quad \langle a^{(N)} | \alpha \rangle^*]$$

Operators: Insert Identity twice

$$\begin{aligned} X &= 1 \times 1 = \left(\sum_{a''} |a''\rangle \langle a''| \right) X \left(\sum_{a'} |a'\rangle \langle a'| \right) \\ &= \sum_{a''} \sum_{a'} |a''\rangle \underbrace{\langle a'' | X | a' \rangle}_{\text{"Matrix Element"}} \langle a' | \end{aligned}$$

$$\text{i.e. } X \doteq \underline{X} = \begin{bmatrix} \langle a^{(1)} | X | a^{(1)} \rangle & \langle a^{(1)} | X | a^{(2)} \rangle & \dots & \langle a^{(1)} | X | a^{(N)} \rangle \\ \langle a^{(2)} | X | a^{(1)} \rangle & \langle a^{(2)} | X | a^{(2)} \rangle & \dots & \dots \\ \vdots & \vdots & \ddots & \vdots \end{bmatrix}$$

Operators are represented by a matrix!

NOTE: $X^\dagger = \sum_{a''} \sum_{a'} |a'\rangle \langle a'' | X^\dagger | a' \rangle \langle a'' |$

$$= \sum_{a''} \sum_{a'} \sum_{a'''} |a'\rangle \langle a' | X | a''' \rangle^* \langle a'' |$$

Hence we
matrix
adjoint!

Spin -1/2 Example

$$S_y = \frac{\hbar}{2} \left[|+\hat{z}\rangle (-i) \langle -\hat{z}| + |-\hat{z}\rangle (i) \langle +\hat{z}| \right]$$

$$\equiv \frac{\hbar}{2} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} = \underline{S}_y = \frac{\hbar}{2} \sigma_y \quad \text{"Pauli Matrices"}$$

NOTE: $\underline{S}_y = \underline{S}_y^\dagger$ i.e. Hermitian Matrix

Operators on Kets

Consider $X|\alpha\rangle = |\beta\rangle \Rightarrow \langle a' | X | \alpha \rangle = \langle a' | \beta \rangle$

~~$\Leftrightarrow \sum_{a''} \sum_{a'} \langle a'' | \alpha \rangle \langle a' | X | a'' \rangle \langle a' | \alpha \rangle = \langle a' | \beta \rangle$~~

$\Leftrightarrow \sum_{a''} \langle a' | X | a'' \rangle \langle a'' | \alpha \rangle = \langle a' | \beta \rangle$

matrix multiplication!!

i.e.
$$\begin{bmatrix} \langle a^{(1)} | X | a^{(1)} \rangle & \langle a^{(1)} | X | a^{(2)} \rangle & \dots \\ \langle a^{(2)} | X | a^{(1)} \rangle & \langle a^{(2)} | X | a^{(2)} \rangle & \dots \\ \vdots & \vdots & \ddots \end{bmatrix} \begin{bmatrix} \langle a^{(1)} | \alpha \rangle \\ \langle a^{(2)} | \alpha \rangle \\ \vdots \end{bmatrix} = \begin{bmatrix} \langle a^{(1)} | \beta \rangle \\ \langle a^{(2)} | \beta \rangle \\ \vdots \end{bmatrix}$$

Similar for Bras, Multiplying Operators

\Leftrightarrow We can do all of our "work" using matrix algebra!

[For Discrete Bases !!]

the Eigenvalue Problem

Suppose you want to "solve" the problem

$$B|b'\rangle = b'|b'\rangle$$

for eigenvalues b' and eigenstates $|b'\rangle$ in terms of some other states a' and $|a'\rangle$.

$$\Leftrightarrow \langle a' | B | b' \rangle = b' \langle a' | b' \rangle$$

$$\sum_{a''} \langle a' | B | a'' \rangle \langle a'' | b' \rangle = b' \langle a' | b' \rangle$$

i.e. $\underline{B} \underline{b} = b' \underline{b}$ Standard Problem!

Example: $S_y |y'\rangle = y' |y'\rangle$ for $s = \frac{1}{2}$

$$\Leftrightarrow \underline{S}_y \underline{y} = y' \underline{y}$$

Found $\underline{S}_y = \frac{\hbar}{2} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \Rightarrow$ handy to write $y' = \frac{\hbar}{2} \lambda$

$$\Leftrightarrow \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = \lambda \begin{bmatrix} y_1 \\ y_2 \end{bmatrix}$$

to find λ :

$$\det \begin{bmatrix} -\lambda & -i \\ i & -\lambda \end{bmatrix} = (-\lambda)^2 - (i)(-i) = \lambda^2 - 1 = 0$$

$$\Leftrightarrow \lambda = \pm 1$$

so $y' = \pm \frac{\hbar}{2}$ CORRECT!

Now find eigenvectors.

For $\lambda = +1$: $\begin{bmatrix} -1 & -i \\ i & -1 \end{bmatrix} \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = \begin{bmatrix} y_1 \\ y_2 \end{bmatrix}$

i.e. $-iy_2 = y_1$ $iy_1 = y_2$ (same)

$\Rightarrow |+\hat{y}\rangle = \frac{1}{\sqrt{2}} |+\hat{z}\rangle + \frac{i}{\sqrt{2}} |-\hat{z}\rangle$

For $\lambda = -1$:

$\begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} = - \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} \Rightarrow \begin{matrix} -iy_2 = -y_1 \\ iy_1 = -y_2 \end{matrix}$

$\Rightarrow |-\hat{y}\rangle = \frac{1}{\sqrt{2}} |+\hat{z}\rangle - \frac{i}{\sqrt{2}} |-\hat{z}\rangle$

NOTE: this is a consistent choice!

i.e. we used these defs of $|\pm\hat{y}\rangle$ to build S_y !

LATER: the physical meaning of angular momentum.