

- * Via Zoom from APS 2024
- * Homework posted
- * New syllabus starts today

Quick Review: Wave Mechanics in 1D

Position operator $x \quad x|x\rangle = x|x\rangle$ "operator \rightarrow eigenvalue"

Translation operator $T(a)|x\rangle = |x+a\rangle$

Normalization $\langle x'|x\rangle = \delta(x-x')$

$$\Leftrightarrow 1 = \int dx |x\rangle \langle x|$$

$T(dx) = 1 - \frac{i}{\hbar} dk p$ "Weil's Trick" $\Rightarrow T(a) = e^{-ipa/\hbar}$
 $p = \text{"momentum"}$ Why?? Well we found that

$$H = \frac{1}{2m} p^2 + V(x) \Rightarrow \frac{d}{dt} \langle x \rangle = \frac{-p}{m} \quad \text{and} \quad \frac{d\langle p \rangle}{dt} = \left\langle -\frac{dV}{dx} \right\rangle \text{ ok}$$

$$\text{Define } \psi_x(x) = \langle x|\alpha \rangle \Rightarrow \langle x|p|\alpha \rangle = \frac{\hbar}{i} \frac{\partial \psi_x}{\partial x}$$

$$\Leftrightarrow H|\psi\rangle = E|\psi\rangle \text{ becomes } -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V(x)\psi(x) = E\psi(x)$$

"Schrödinger Equation"

- Also
- For "Bound states" have $\psi(x \rightarrow \pm\infty) = 0$
 - \Leftrightarrow Discrete Energies! (e.g. Box, SHO)
 - $\langle x|p|\psi\rangle = p\langle x|\psi\rangle$ $\Rightarrow \langle x|\psi\rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ixp/\hbar}$
 $\text{and } = \frac{\hbar}{i} \frac{\partial}{\partial x} \langle x|\psi\rangle$

Today: Begin Wave Mechanics in 3D

Position and Momentum in 3D

Position Eigenstate $| \vec{r} \rangle = | x, y, z \rangle$

of operator $\vec{r} = x \hat{x} + y \hat{y} + z \hat{z}$ $\hat{x}, \hat{y}, \hat{z}$ unit vectors

$$\text{i.e. } x | \vec{r} \rangle = x | \vec{r} \rangle \quad y | \vec{r} \rangle = y | \vec{r} \rangle \quad z | \vec{r} \rangle = z | \vec{r} \rangle$$

$\xrightarrow{\text{operator}}$ $\xrightarrow{\text{eigenvalue}}$

Tacit assumption: x, y, z are commuting operators!

Don't assume that \rightarrow "super symmetry" (!)

Normalization

$$\langle \vec{r} | \vec{r}' \rangle = \delta(x-x') \delta(y-y') \delta(z-z')$$

$$= \delta^{(3)}(\vec{r} - \vec{r}')$$

Wave Function $\psi_{\alpha}(\vec{r}) = \langle \vec{r} | \alpha \rangle$

$$\Leftrightarrow 1 = \langle \alpha | \alpha \rangle = \int d^3r \langle \alpha | \vec{r} \rangle \langle \vec{r} | \alpha \rangle = \int d^3r \psi_{\alpha}^*(\vec{r}) \psi_{\alpha}(\vec{r})$$

Translation Operator: Now it has a vector argument

$$T(a_x \hat{x}) | x, y, z \rangle = | x+a_x, y, z \rangle$$

$$T(a_y \hat{y}) | x, y, z \rangle = | x, y+a_y, z \rangle$$

$$T(a_z \hat{z}) | x, y, z \rangle = | x, y, z+a_z \rangle$$

$$\text{i.e. } T(\vec{a}) | \vec{r} \rangle = | \vec{r} + \vec{a} \rangle$$

We clearly have

$$T(a_x \hat{x}) = e^{-i p_x a_x / \hbar}$$

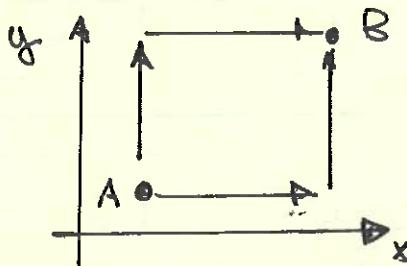
$$T(a_y \hat{y}) = e^{-i p_y a_y / \hbar}$$

$$T(a_z \hat{z}) = e^{-i p_z a_z / \hbar}$$

same argument as for "x"

i.e. $a_x = \lim_{N \rightarrow \infty} \frac{q_k}{N}$ etc...

Important: Translations commute.



i.e. $T(a_y \hat{y}) T(a_x \hat{x})$
 $= T(a_x \hat{x}) T(a_y \hat{y})$

↳ Put $a_x = dx$, $a_y = dy$ and go to second order
 to find $[p_x, p_y] = 0$ Homework!

"Momentum components commute" $\Rightarrow |p_x, p_y, p_z\rangle = |\vec{p}\rangle$ ok!

This also means that we can write...

$$\begin{aligned} T(a_x \hat{x}) T(a_y \hat{y}) T(a_z \hat{z}) &= e^{-i(p_x a_x + p_y a_y + p_z a_z) / \hbar} \\ &= e^{-i \vec{p} \cdot \vec{a} / \hbar} \end{aligned}$$

... because the p_x, p_y, p_z all commute!

Back to 1D for a moment

We showed $(x T(\delta x) - T(\delta x) x) = -\frac{i}{\hbar} \delta x [x, p_x]$

and $(x T(\delta x) - T(\delta x) x) |\alpha\rangle = \delta x |\alpha\rangle \Leftrightarrow x T(\delta x) \rightarrow (x + \delta x)$

$\Leftrightarrow [x, p_x] = i\hbar$ "Canonical commutation Relation"

so it follows directly that

$$[y, p_y] = i\hbar \quad \text{and} \quad [z, p_z] = i\hbar$$

But what about other combinations?

i.e. Do y and $T(\alpha_x \hat{x})$ commute?

$$\begin{aligned} T(\alpha_x \hat{x}) y |\alpha\rangle &= T(\alpha_x \hat{x}) y \int d^3r |x, y, z\rangle \langle x, y, z | \alpha \rangle \\ &= T(\alpha_x \hat{x}) \underbrace{\int d^3r y |x, y, z\rangle \langle x, y, z | \alpha \rangle}_{\substack{\text{operator} \\ \text{eigenvalue}}} \xrightarrow{\text{i.e. } y T(\alpha_x \hat{x}) \text{ has no eigenvalue change!}} \\ &= \int d^3r y T(\alpha_x \hat{x}) |x, y, z\rangle \langle x, y, z | \alpha \rangle = y T(\alpha_x \hat{x}) |\alpha\rangle \end{aligned}$$

yes $\Rightarrow [y, p_x] = 0$ etc...

$$\text{i.e. } \underline{[r_i, p_j]} = i\hbar \delta_{ij} \quad r_1, r_2, r_3 = x, y, z$$

$$p_1, p_2, p_3 = p_x, p_y, p_z$$

Nonrelativistic Eigenstates $|p_x, p_y, p_z\rangle = |\vec{p}\rangle$

$$\text{or } \langle \vec{p} | \vec{p}' \rangle = \delta^{(3)}(\vec{p} - \vec{p}')$$

$$\begin{aligned} \text{and as before } \langle \vec{r} | \vec{p} | \alpha \rangle &= \frac{i}{\hbar} \left(\hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z} \right) \langle \vec{r} | \alpha \rangle \\ &= \frac{i}{\hbar} \vec{\nabla} \underline{\Psi}_{\alpha}(\vec{r}) \end{aligned}$$

$$\text{and } \langle \vec{r} | \vec{p} \rangle = \frac{1}{(2\pi\hbar)^{3/2}} e^{i \vec{p} \cdot \vec{r} / \hbar}$$

Translational Invariance

aka "Translational Symmetry"

Consider a two-particle Hamiltonian:

$$H = \frac{1}{2m_1} \vec{p}_1^2 + \frac{1}{2m_2} \vec{p}_2^2 + V(|\vec{r}_1 - \vec{r}_2|)$$

Only depends on difference!

↳ Expect translational invariance!

Formal: $T_{12}^\dagger H T_{12} = H$ (See homework)

$$\text{w/ } T_{12} = T_1(\vec{q}) T_2(\vec{q}) = e^{-i \vec{P} \cdot \vec{a} / \hbar}$$

where $\vec{P} = \vec{p}_1 + \vec{p}_2$ "total momentum"

Find $[\vec{P}, H] = 0$

But $\frac{d}{dt} \langle \vec{P} \rangle = \frac{i}{\hbar} \langle [H, \vec{P}] \rangle = 0$ || etc...

↳ "Total Momentum is conserved"

Recall: $d\langle A \rangle / dt$
 $= d/dt \langle \alpha_i | A | \alpha_i \rangle$

We saw the same result in

Analytical Mechanics!

"Conservation Laws are Fundamental"

↳ Much of this physics follows directly from classical mechanics

e.g. Relative and CM Coordinates

Better operators for two-body problems

$$\vec{r} \equiv \vec{r}_1 - \vec{r}_2 \quad \text{"relative"}$$

$$\vec{R} \equiv (m_1 \vec{r}_1 + m_2 \vec{r}_2) / (m_1 + m_2) \quad \text{"CM"}$$

↓ Might not have
true for this,
but that's ok.

NOTE $[\vec{r}_i, \vec{p}_j] = 0$

$$\text{e.g. } [x_1, -x_2, \vec{p}_1 + \vec{p}_2] = [x_1, \vec{p}_1] - [x_2, \vec{p}_2] = 0$$

Expected from translational invariance!

Also $[\vec{R}_i, \vec{p}_j] = i\hbar \delta_{ij}$ (Work it out)

↳ Define relative momentum $\vec{p} = \frac{m_2 \vec{p}_1 - m_1 \vec{p}_2}{m_1 + m_2}$

↳ $[\vec{r}_i, \vec{p}_j] = i\hbar \delta_{ij}$ (Work it out)

But momentum thus:

↳ we can treat two-body systems just using the relative coordinate and reduced mass.
(As Expected).

Upshot: We will now consider interactions

$$H = \frac{1}{2m} \vec{p}^2 + V(\underline{\vec{r}}) \quad \text{"Relative } \vec{r} \text{ and } \vec{p}$$

in Two-Body System"

$$m = \frac{m_1 m_2}{m_1 + m_2} \quad \text{"Reduced Mass"}$$

"Central Potential"

i.e. spherically symmetric