

Reproducibility in Computational and Experimental Mathematics

Dec. 10–14, 2012

Abstract

We present a simple method to solve spherical harmonics moment systems, such as the the time-dependent P_N and SP_N equations, of radiative transfer. The method, which works for arbitrary moment order **N**, makes use of the specific coupling between the moments in the P_N equations. This coupling naturally induces staggered grids in space and time, which in turn give rise to a canonical, second-order accurate finite difference scheme. While the scheme does not possess TVD or realizability limiters, its simplicity allows for a very efficient implementation in MATLAB. We present several test cases, some of which demonstrate that the code solves problems with ten million degrees of freedom in space, angle, and time within a few seconds. The code for the numerical scheme, called StaRMAP (Staggered grid Radiation Moment **Approximation**), along with files for all presented test cases, can be downloaded so that all results can be reproduced by the reader.

Radiative Transfer Equation for Photons in Medium

$\partial_t \psi(t, \mathbf{x}, \Omega) + \Omega \cdot \nabla_{\mathbf{x}} \psi(t, \mathbf{x}, \Omega) + \Sigma_t(t, \mathbf{x}) \psi(t, \mathbf{x}, \Omega)$	<i>,</i>
$= \int_{S^2} \Sigma_s(t, x, \Omega \cdot \Omega') \psi(t, x, \Omega') \mathrm{d}\Omega' + q(t, x, \Omega)$	(1

Photon density ψ ; absorption cross section Σ_a ; scattering kernel Σ_s ; total cross section $\Sigma_t = \Sigma_{s0} + \Sigma_a$; source q.

Spherical Harmonic Moment Methods

The P_N method (cf. [1]) conducts a Fourier expansion (spectral discretization) in the angular variable Ω ; reduces high dimensionality; yields system of macroscopic PDE.

Efficient numerics recent subject of interest [5, 2, 1].

Advantages over more direct discretizations: rotational invariance, no ray effect (cf. [3]).

Drawback: Gibbs phenomena, i.e., spurious oscillations

Also considered: Simplified P_N equations (SP_N) [4].

Gap: No P_N solver for general N available.

Moment system is hyperbolic balance law

$\partial_t \vec{u} + M_x \cdot \partial_x \vec{u} + M_y \cdot \partial_y \vec{u} + C \cdot \vec{u} = \vec{q} ,$ (2)

where the matrices M_x , M_y , C possess very specific patterns of their nonzero entries that admit the placement of the components of the solution vector $\vec{u}(x, y, t)$ on staggered grids.

	 grid 11 grid 21 grid 12 grid 22 perio- dic exten- sion extra- polation
--	---

Staggered grid of 5×3 cells; periodic b.c. in **x**; extrapolation b.c. in **y**; solution grid points (black boundaries), periodic extension points (blue), and extrapolation ghost points (red).

Features of This Project [6]

Numerical Method

 $\partial_{\tau} u_k(\mathbf{x}, \mathbf{y}, \tau) + \bar{c}_k(\mathbf{x}, \mathbf{y}) u_k(\mathbf{x}, \mathbf{y}, \tau) = \bar{r}_k(\mathbf{x}, \mathbf{y}) \quad (3)$ for $\tau \in [t, t + \Delta t]$. Evaluate $\overline{c}_k(x, y) = c_k(x, y, t + \frac{1}{2}\Delta t)$ and $\overline{r}_k(x, y) = r_k(x, y, t + \frac{1}{2}\Delta t)$. Exact solution of (3) is $u_k(x, y, t + \Delta t) = u_k(x, y, t) + \Delta t (\bar{r}_k(x, y))$ $-\bar{c}_k(x,y)u_k(x,y,t)) E(-\bar{c}_k(x,y)\Delta t),$ where $E(c) = \frac{\exp(c)-1}{c}$.

Accuracy: 2nd order due to local symmetries. Stability: Proof in [6].



Project Files Used by MMS Verification

Example file creates new file which calls P_N matrices constr. and then solver file

MMS source is computed via MATLAB's symbolic toolbox.

StaRMAP

A Second Order Staggered Grid Method for Spherical Harmonics Moment Equations of Radiative Transfer

• Specific MATLAB files that encode P_N and SP_N matrices. Efficient solver file:

▶ Place (based on matrices M_x and M_y) solution components

automatically on appropriate staggered grids.

Store solution components as 2d arrays (= matrices). Compute FD stencils by shifting. Very fast in MATLAB.

► Solver file employs special structure of parameters (e.g., isotropy, time-independence) to compute much faster.

Examples files (few lines of MATLAB code) that call the other files. User hardly ever modifies non-example files. Solver file has defaults for every problem parameter and function, so an example file essentially is a prescription of how one deviates from the defaults.

Central differencing in space:

 $\partial_x w_{11} \& \partial_y w_{22}$ live on 21 grid; $\partial_x w_{22} \& \partial_y w_{11}$ live on 12 grid; $\partial_x w_{21} \& \partial_y w_{12}$ live on 11 grid; $\partial_x w_{12} \& \partial_y w_{21}$ live on 22 grid. Even components: on grids 11 and 22; odd components: on grids 21 and 12.

Bootstrapping in time: Update even (odd) components from t to $t + \Delta t$, assuming that odd (even) components are constant. Thus system decouples into scalar ODEs

> starmap_create_mms.m starmap_ex_mms_auto.m starmap_closure_pn.m starmap_solver.m



Project Files Used by Checkerboard Test Case

Example file starmap_ex_checkerboard.m calls P_N matrices construction starmap_closure_pn.m and then solver file starmap_solver.m

StaRMAP Example File for Checkerboard Test Case

%=====================================
<pre>% Problem Parameters %====================================</pre>
<pre>par = struct('name','Checkerboard Test', % name of 'closure','P', % type of closure (ca 'n_mom',5, % order of moment approxi 'sigma_a',@sigma_a, % absorption coe 'sigma_s0',@sigma_s0, % isotropic sc 'source',@source, % source term (des 'ax',[0 7 0 7], % coordinates of con 'n',[250 250], % numbers of grid ces 'bc',[1 1], % type of boundary cond 't_plot',linspace(0,3.2,51), % output 'output',@output % problem-specific);</pre>
<pre>% Moment System Setup and Solver Execut</pre>
<pre>%====================================</pre>
<pre>% Problem Specific Functions</pre>
<pre>%====================================</pre>
<pre>% Absorption coefficient. cx = ceil(x); cy = ceil(y); g = (ceil((cx+cy)/2)*2==(cx+cy)).*(1<c; f = (1-g)*0+g*10;</c; </pre>
<pre>% Absorption coefficient. cx = ceil(x); cy = ceil(y); g = (ceil((cx+cy)/2)*2==(cx+cy)).*(1<c: f = (1-g)*0+g*10; function f = sigma_s0(x,y) % Isotropic scattering coefficient. cx = ceil(x); cy = ceil(y); g = (ceil((cx+cy)/2)*2==(cx+cy)).*(1<c: f = (1-g)*1+g*0;</c: </c: </pre>
<pre>% Absorption coefficient. cx = ceil(x); cy = ceil(y); g = (ceil((cx+cy)/2)*2==(cx+cy)).*(1<c: f = (1-g)*0+g*10; function f = sigma_s0(x,y) % Isotropic scattering coefficient. cx = ceil(x); cy = ceil(y); g = (ceil((cx+cy)/2)*2==(cx+cy)).*(1<c: f = (1-g)*1+g*0; function f = source(x,y) % Radiation source (only for zeroth mon f = 3<x&x<4&3<y&y<4;< pre=""></x&x<4&3<y&y<4;<></c: </c: </pre>

Benjamin Seibold & Martin Frank Temple University & RWTH Aachen University





ic color scale. ght, caxis([-7 0])2f',par.name,par.closure,...



