

Interface Tracking

Many two-phase fluid flow simulations (e.g. immersed boundary method [6], ghost fluid method [2]) use fractional steps, i.e. in each time step:

1. Update velocity field $\vec{v}(\vec{x})$ by a Navier-Stokes step;
2. Move interface by the velocity field $\vec{v}(\vec{x})$.

Fundamental problem: Represent and move the interface accurately.

Level Set Approach [5]

- Represent interface as the zero contour of a level set function $\phi : \mathbb{R}^p \rightarrow \mathbb{R}$.
- Moving interface with velocity field \vec{v} translates to evolving ϕ by the PDE

$$\phi_t + \vec{v} \cdot \nabla \phi = 0. \quad (1)$$

- Solve (1) by high accuracy schemes (e.g. WENO [3]).
- Preserve $|\nabla \phi(\mathbf{x})| = 1$ by reinitialization [9] equation $\phi_\tau = \text{sign}(\phi_0)(1 - |\nabla \phi|)$.
- Reconstruct interface by bi/tri-linear interpolation on grid cells.

Advantages:

- Can define ϕ on regular grid; uniform resolution.
- Simple in 3D; automatic treatment of topology changes.
- More potential due to knowledge of *all* contours.
- Can obtain normals $\vec{n} = \frac{\nabla \phi}{|\nabla \phi|}$ and curvature $\kappa = \nabla \cdot \vec{n}$ from ϕ .

Problems and difficulties:

- Small structures vanish once below grid resolution.
- Wide WENO stencils are difficult to deal with near boundaries and with AMR.
- Inaccuracies with curvature approximation near boundaries and with reinitialization.

Goal: Address these problems in a simple, Eulerian fashion.

Idea: Augment the level set function by gradient information.

Gradient-Augmented Level Set Method and Advection Scheme

Philosophy:

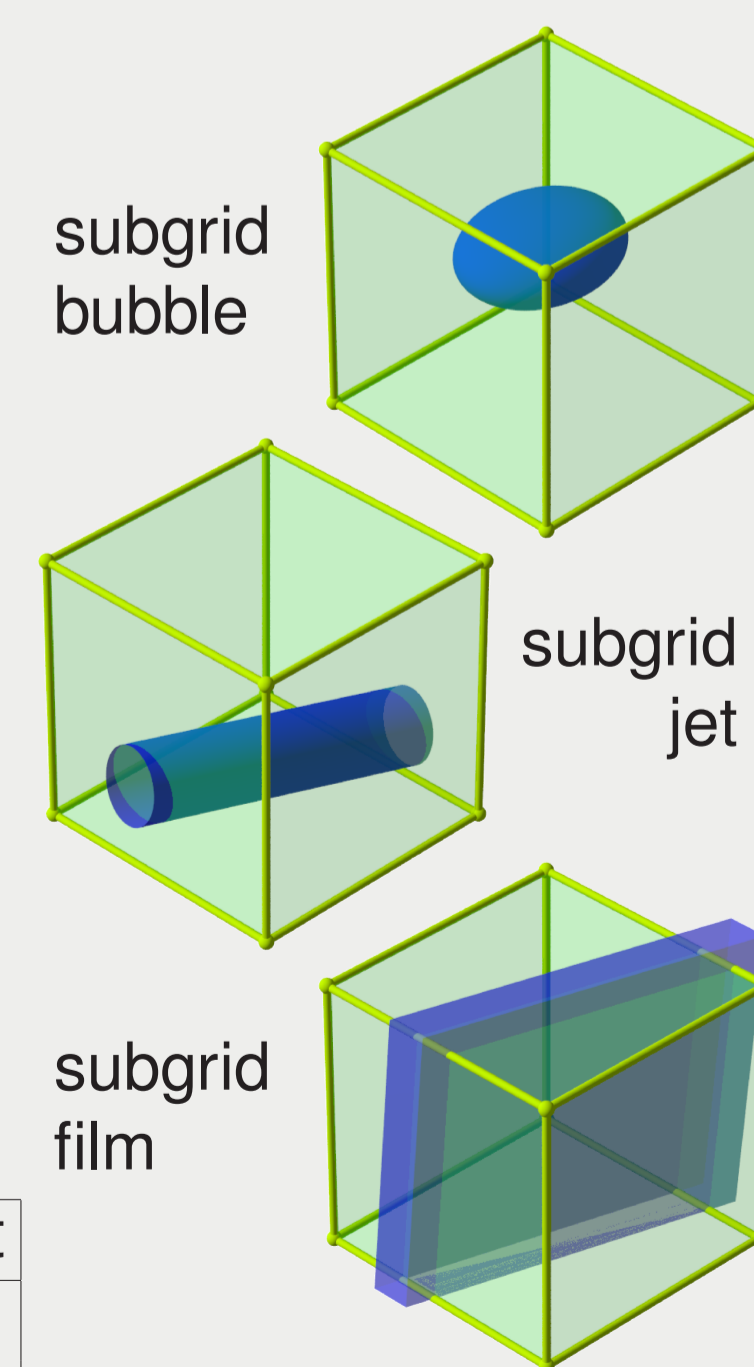
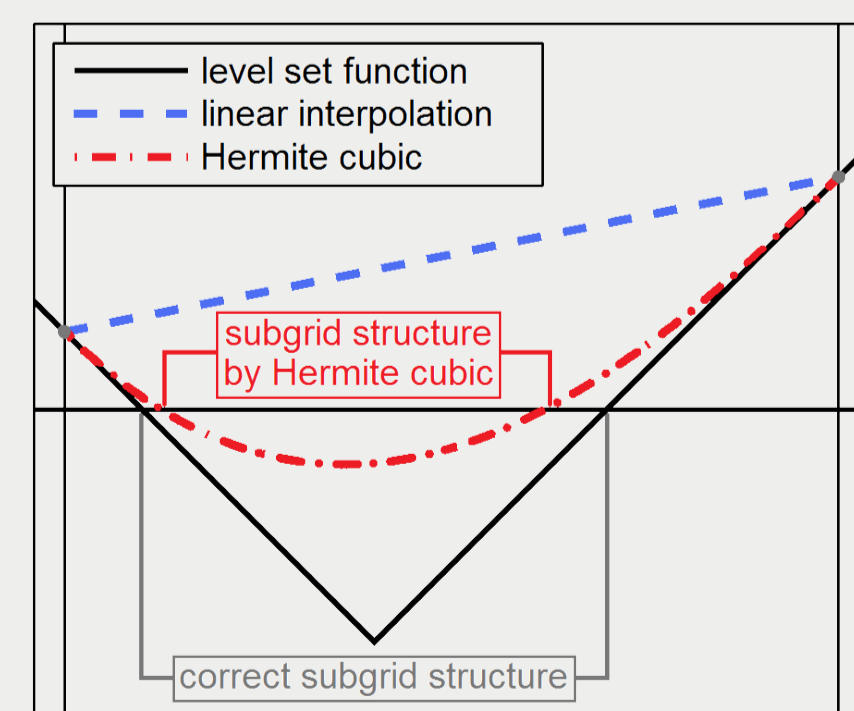
Store function values ϕ and gradients $\vec{\psi} = \nabla \phi$ as independent quantities. Update both in a coupled fashion, by a scheme that uses characteristics and Hermite interpolation [4].

Research in the past:

- Courant et al.: use characteristics and interpolation, but no gradients \rightarrow CIR method [1]
- van Leer: use gradients for higher order, but reconstruct from ϕ [11]
- Takekawa et al.: idea to track gradients as independent quantity \rightarrow CIP method [10]
- Raessi et al.: advect normal vectors, but in a decoupled fashion [7]

Motivation:

If function values ϕ and gradients $\nabla \phi$ are known at grid points, then subgrid structures can be represented.



Hermite interpolation:

Here: p -cubic; simple tensor-product on rectangular cell.

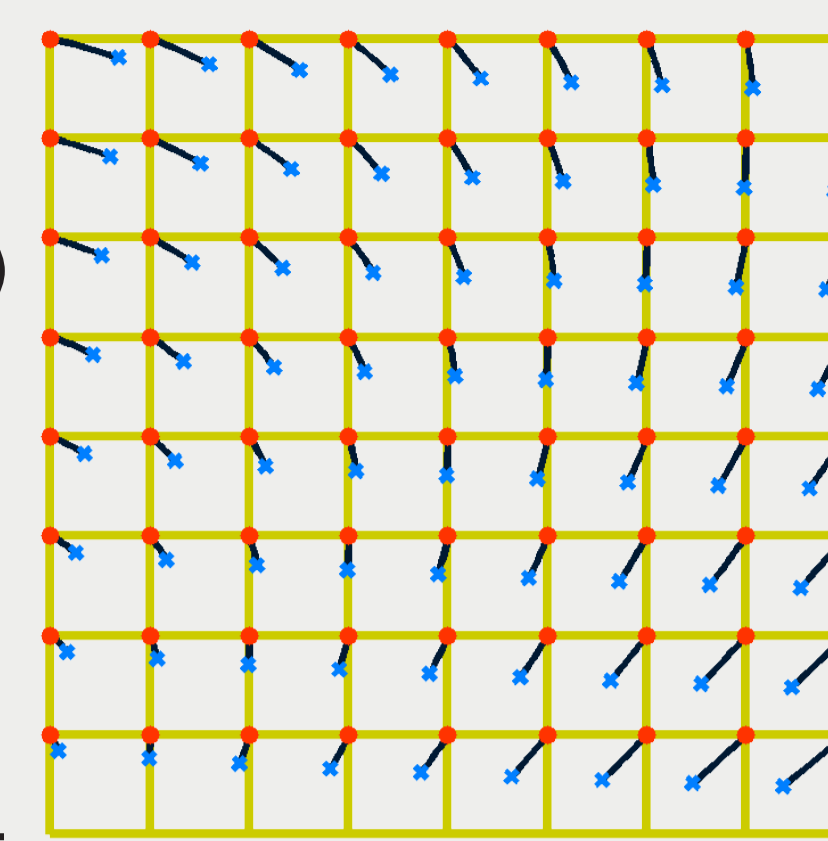
p	given on 2^p vertices	\rightarrow construct approx. to	\rightarrow interpolant
1	ϕ, ϕ_x	$\rightarrow \phi_{xy}$	\rightarrow cubic
2	ϕ, ϕ_x, ϕ_y	$\rightarrow \phi_{xy}, \phi_{xz}, \phi_{yz}, \phi_{xyz}$	\rightarrow bi-cubic
3	$\phi, \phi_x, \phi_y, \phi_z$	$\rightarrow \phi_{xy}, \phi_{xz}, \phi_{yz}, \phi_{xyz}$	\rightarrow tri-cubic

Thm.: The p -cubic approximates smooth functions with 4^{th} order accuracy. Gradients are 3^{rd} order accurate. Curvature can be obtained everywhere with 2^{nd} order accuracy.

Semi-Lagrangian Advection Scheme

Characteristic form of (1)

$$\begin{cases} \dot{\vec{x}} = \vec{v}(\vec{x}) \\ \dot{\phi} = 0 \\ \dot{\vec{\psi}} = -\nabla \vec{v} \cdot \vec{\psi} \end{cases} \quad (2)$$



1. For each grid point, solve $\dot{\vec{x}} = \vec{v}(\vec{x})$ backwards to find starting point of characteristic \vec{x} (e.g. by one RK3 step).
2. Obtain $\phi(\vec{x})$ and $\vec{\psi}(\vec{x})$ from Hermite interpolant.
3. Solve $\dot{\phi} = 0$ (trivial) and $\dot{\vec{\psi}} = -\nabla \vec{v} \cdot \vec{\psi}$ forward (e.g. RK2).

Numerical Analysis of Gradient-Augmented Schemes

Introduce concept of “superconsistency” to analyze scheme in function space setting:

- Rewrite (1) as $\phi_t = L\phi$, where $L = -\vec{v} \cdot \nabla$.
- Exact advection operator $S(t) = e^{tL}$.
- Approximate operator $A(t) \approx S(t)$, where $A(k)$ is one Runge-Kutta step of $\dot{\vec{x}} = \vec{v}(\vec{x})$.

Def.: A gradient-augmented scheme is called **superconsistent**, if the numerical approximation of $\dot{\vec{\psi}} = -\nabla \vec{v} \cdot \vec{\psi}$ equals $\nabla(A(t)\phi)$.

Update rule for gradients is inherited from ODE solver for characteristic curves.

Projection operator P in function space: (1) evaluate ϕ and $\nabla \phi$ on grid points; (2) construct new function by Hermite interpolation on each cell. Thus one step of scheme:

$$\underbrace{PA(k)}_{\text{num. scheme}} \approx \underbrace{PS(k)}_{\text{exact evolution}} \approx \underbrace{S(k)}$$

Example: superconsistent Shu-Osher [8] scheme

$$\begin{aligned} \vec{x}_1 &= \vec{x} - k \vec{v}(\vec{x}, t+k) \\ \nabla \vec{x}_1 &= I - k \nabla \vec{v}(\vec{x}, t+k) \\ \vec{x}_2 &= \vec{x} - k(\frac{1}{3} \nabla \vec{v}(\vec{x}, t+k) + \frac{2}{3} \nabla \vec{v}(\vec{x}_1, t)) \\ \nabla \vec{x}_2 &= I - k(\frac{1}{3} \nabla \nabla \vec{v}(\vec{x}, t+k) + \frac{2}{3} \nabla \nabla \vec{v}(\vec{x}_1, t)) \\ \vec{x} &= \vec{x} - k(\frac{1}{6} \nabla \vec{v}(\vec{x}, t+k) + \frac{4}{6} \nabla \vec{v}(\vec{x}_1, t) + \frac{1}{6} \nabla \vec{v}(\vec{x}_2, t + \frac{1}{2}k)) \\ \nabla \vec{x} &= I - k(\frac{1}{6} \nabla \nabla \vec{v}(\vec{x}, t+k) + \frac{4}{6} \nabla \nabla \vec{v}(\vec{x}_1, t) + \frac{1}{6} \nabla \nabla \vec{v}(\vec{x}_2, t + \frac{1}{2}k)) \\ \phi(\vec{x}, t+k) &= \phi(\vec{x}, t) \\ \vec{\psi}(\vec{x}, t+k) &= \nabla \vec{x} \cdot \vec{\psi}(\vec{x}, t) \end{aligned}$$

Accuracy:

Projection P fourth order accurate. Use locally fourth order accurate ODE solver for $\dot{\vec{x}} = \vec{v}(\vec{x})$. Then the scheme is locally fourth order accurate, globally third order.

Stability:

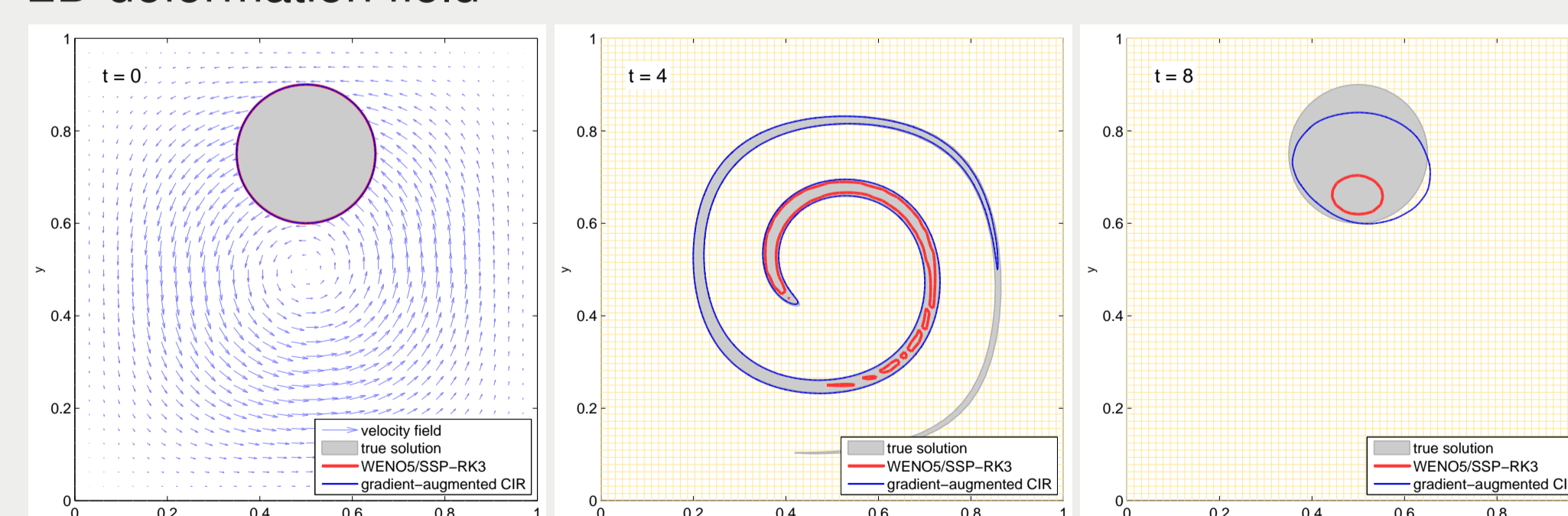
1D constant coefficients: Hermite cubic minimizes $F(u) = \int u_{xx}^2 dx$; hence $F(PA(k)\phi) \leq F(A(k)\phi) = F(S(k)\phi) = F(\phi)$.

Convergence:

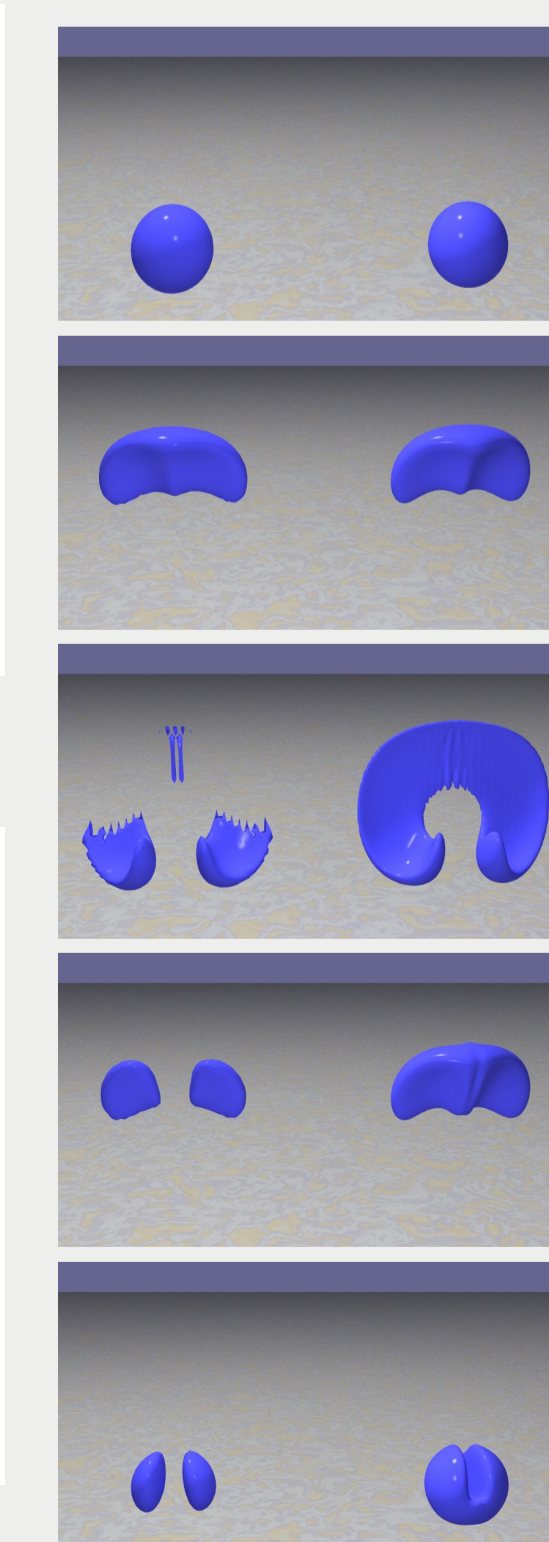
Linear scheme, thus convergence due to the Lax equivalence theorem.

Numerical Benchmark Tests

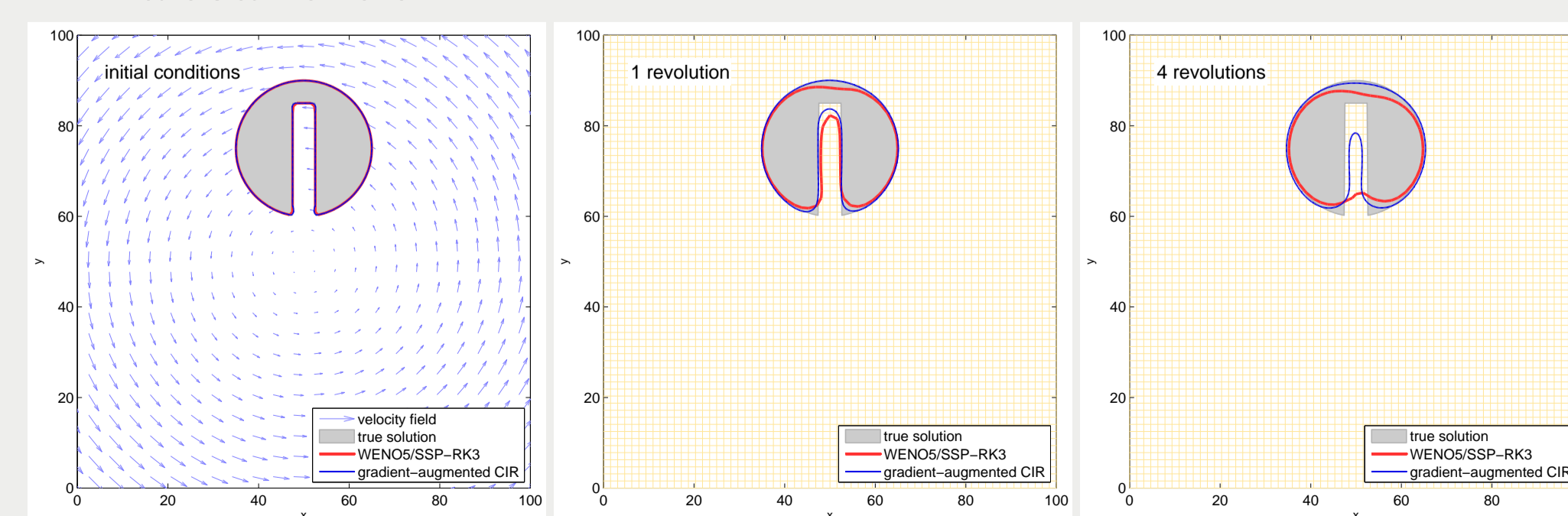
2D deformation field



3D deformation



2D Zalesak circle

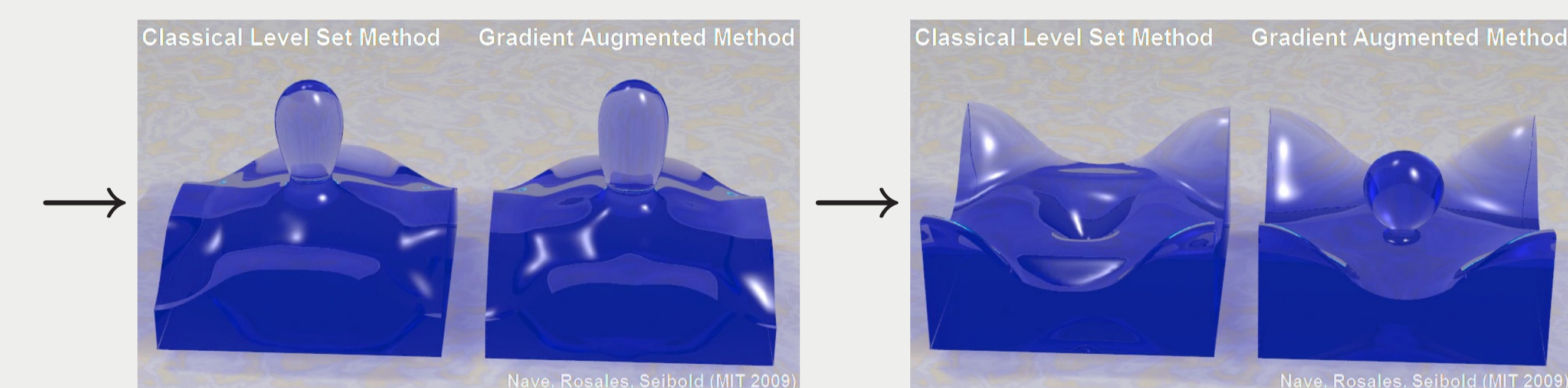


Loss of Volume

resolution	50 × 50 × 50	100 × 100 × 100	150 × 150 × 150	200 × 200 × 200
GA-CIR	-5.4%	-3.3%	-1.6%	-0.9%
WENO5-TVD3	-54.5%	-15.8%	-9.6%	-6.5%
WENO5 reinit.	-35.4%	-11.2%	-7.1%	-5.0%

Evolution of a cube in a 3D deformation field. Lowest volume loss with the new approach.

Fluid Flow Simulation (Using Ghost Fluid Method [2])



The gradient-augmented approach for interface evolution yields a more accurate resolution of partial coalescence events. The ability to resolve small structures pays off.

Conclusions and Outlook

Advantages of gradient-augmented level set approaches:

- Representation of subgrid structures.
- Optimal locality: each point uses only information in a single cell. Great benefit for AMR.
- Accurate approximation to \vec{n} and κ everywhere (directly from interpolation).
- Characteristics yield a natural treatment of boundary conditions.
- Robust implementation: use interpolation for everything.

Current research:

- Analysis: extend concept of TVD to gradient-augmented setting.
- Extension: use other Hermite interpolations than p -cubics.
- Extension: track higher derivatives.
- Extension: nonlinear Hamilton-Jacobi equations (such as the reinitialization equation).

References

[1] R. Courant, E. Isaacson, M. Rees, *On the solution of nonlinear hyperbolic differential equations by finite differences*, Comm. Pure Appl. Math. 5, pp. 243–255, 1952.

[2] R. Fedkiw, T. Aslam, B. Merriman, S. Osher, *A non-oscillatory Eulerian approach to interfaces in multimaterial flows (the ghost fluid method)*, J. Comput. Phys. 152, pp. 457–492, 1999.

[3] X.-D. Liu, S. Osher, T. Chan, *Weighted essentially non-oscillatory schemes*, J. Comput. Phys. 115, pp. 200–212, 1994.

[4] J. Nave, R. R. Rosales, B. Seibold, *A gradient-augmented level set method with an optimally local, coherent advection scheme*, J. Comput. Phys. 229, pp. 3802–3827, 2010. arXiv:0905.3409 [math-ph]

[5] S. Osher, J. A. Sethian, *Fronts propagating with curvature-dependent speed: Algorithms based on Hamilton–Jacobi formulations*, J. Comput. Phys. 79, pp. 12–49, 1988.

[6] C. S. Peskin, *Flow patterns around heart valves: A numerical method*, J. Comput. Phys. 10, pp. 252–271, 1972.

[7] M. Raessi, J. Mostaghimi, M. Bussmann, *Advecting normal vectors: A new method for calculating interface normals and curvatures when modeling two-phase flows*, J. Comput. Phys. 226(1), pp. 774–797, 2007.

[8] C.-W. Shu, S. Osher, *Efficient implementation of essentially non-oscillatory shock-capturing schemes*, J. Comput. Phys. 77, pp. 439–471, 1988.

[9] M. Sussman, P. Smereka, S. Osher, *A level set approach for computing solutions to incompressible two-phase flow*, J. Comput. Phys. 114(1), pp. 146–159, 1994.

[10] H. Takekawa, A. Nishiguchi, T. Yabe, *Cubic interpolated pseudoparticle method (CIP) for solving hyperbolic type equations*, J. Comput. Phys. 61, pp. 261–268, 1985.

[11] B. van Leer, *Towards the ultimate conservative difference scheme V. A second-order sequel to Godunov’s method*, J. Comput. Phys. 32, pp. 101–136, 1979.